

Baryon acoustic oscillations

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Baryons

- ◆ We have neglected the **baryons**: we suppose that dark matter is the only matter component.
- ◆ Link with observation: we assume that the distribution of baryons **follow** that of the dark matter.
- ◆ The baryons **do not contribute** to the gravitational potential: they just fall into it.
- ◆ **Successful** approach: the baryons constitute only 15 percent of the total matter.
- ◆ Missing point: **baryon acoustic oscillations**.

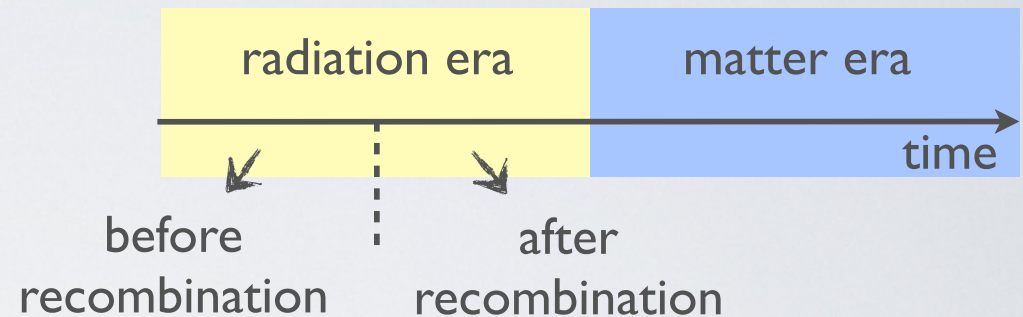
Oscillations

Difference between baryons and dark matter:
interaction with **photons**.

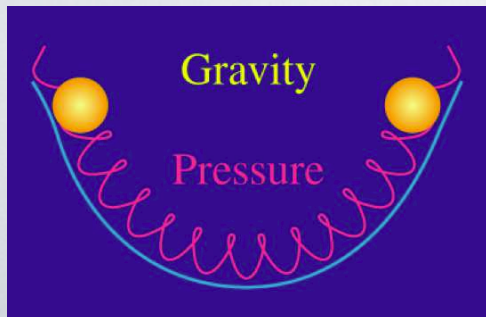
◆ Dark matter evolution: two stages



◆ Baryons evolution: three stages



Before recombination: photons and baryons are tightly **coupled**.
They behave as a **single fluid**, governed by two opposite forces.



- ◆ gravitational attraction
- ◆ radiation pressure



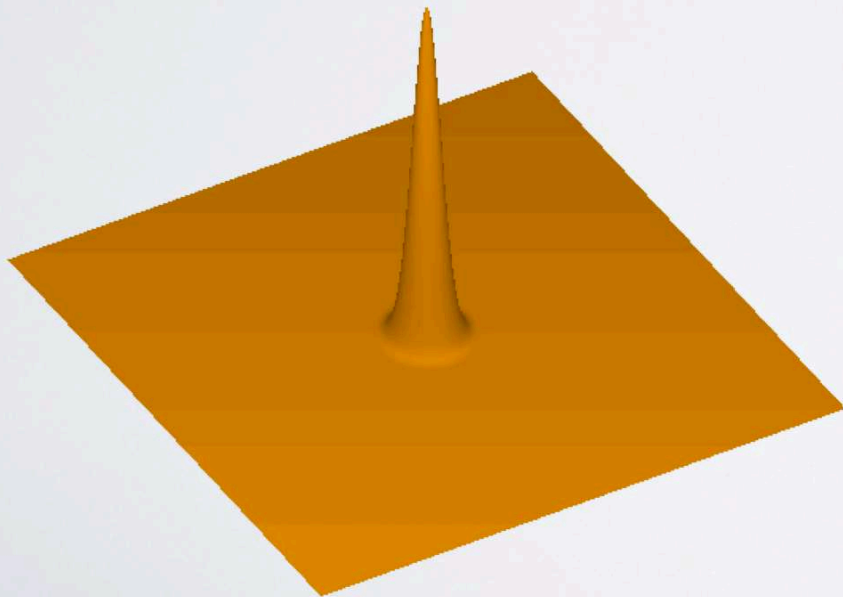
send a **sound wave**
through the universe

Sound wave

- ◆ At **recombination**: electrons and protons combine to form neutral hydrogen. The interaction with the photons stops.
- ◆ This process leaves an impact on the matter **power spectrum**.

In real space

Toy model: an over-density at a single point in the universe.



The fluid propagate until recombination:

- ◆ the photons escape
- ◆ the baryons stay

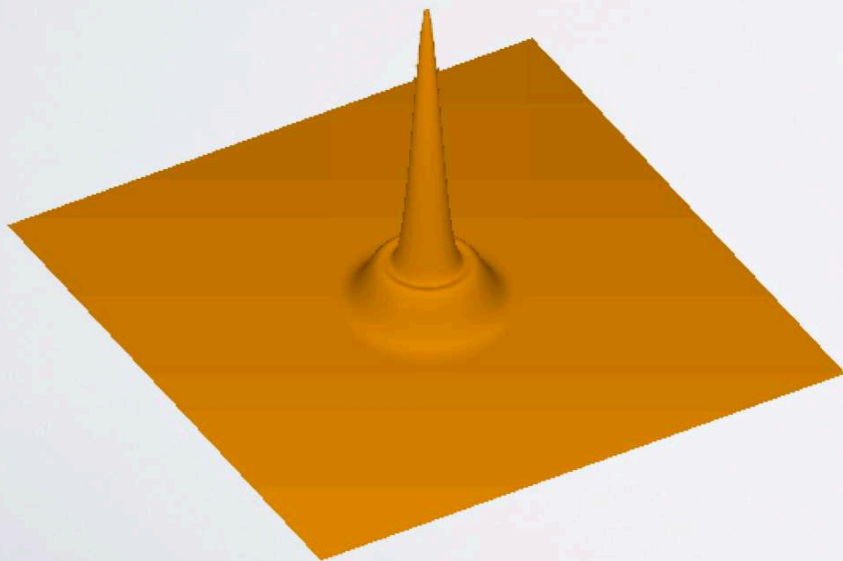
Shell of **baryons** at distance 150 Mpc

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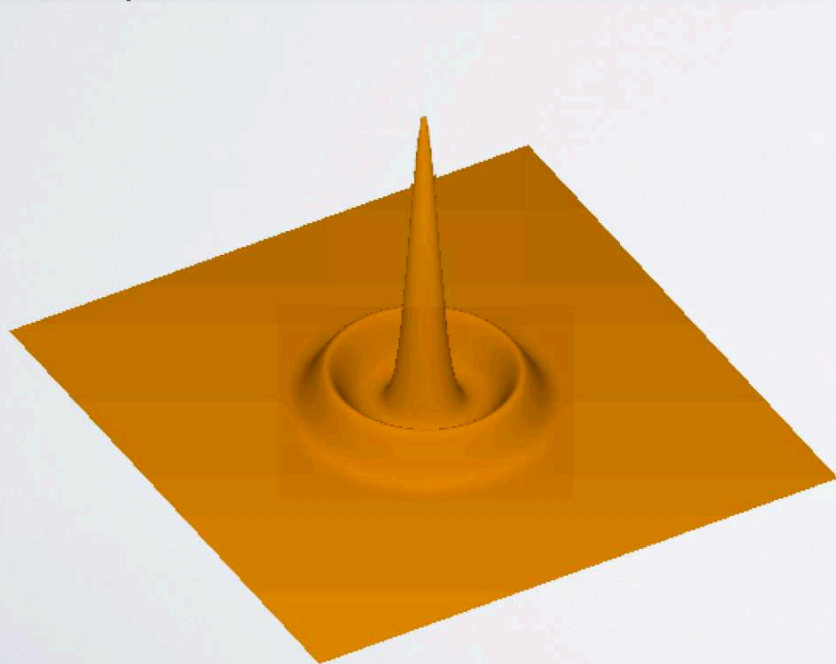
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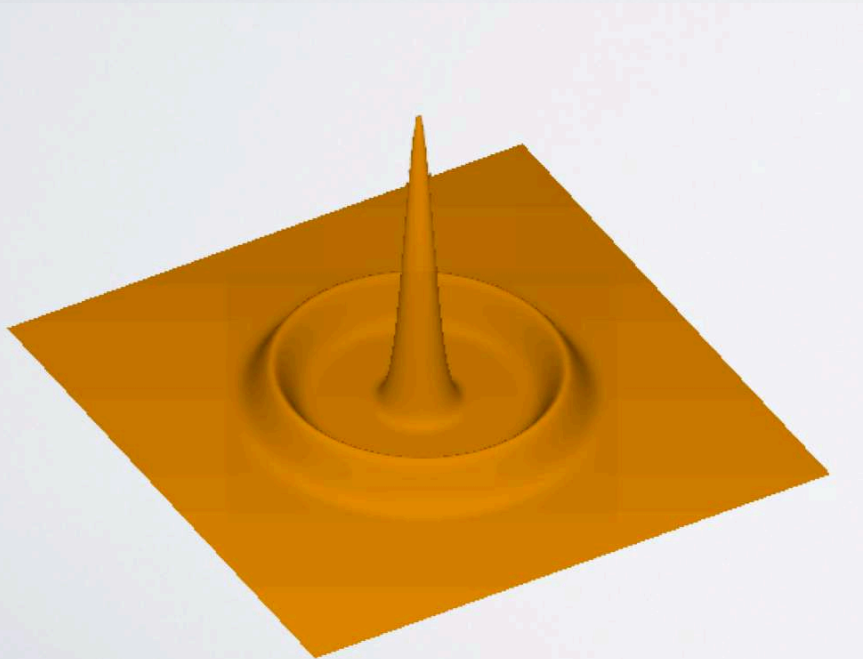
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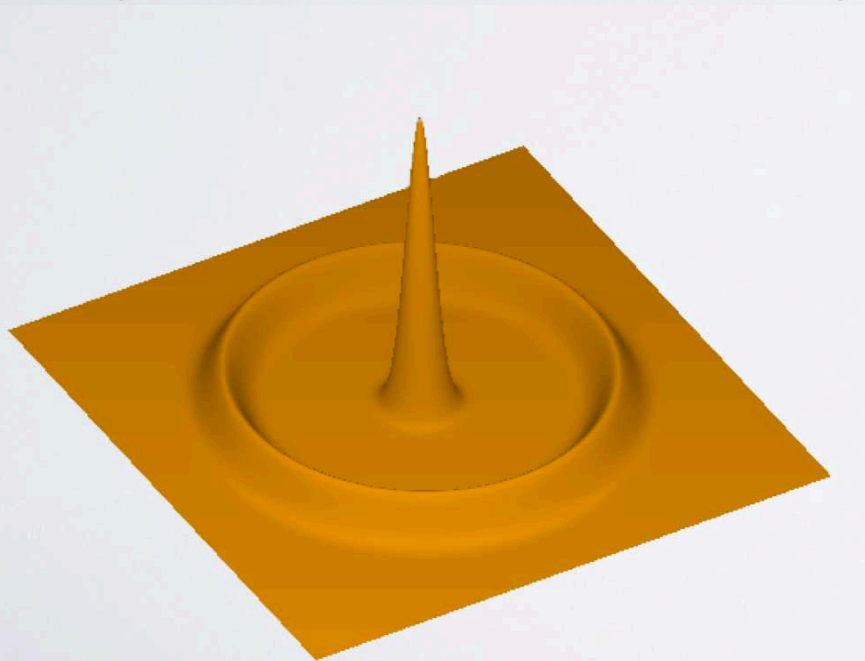
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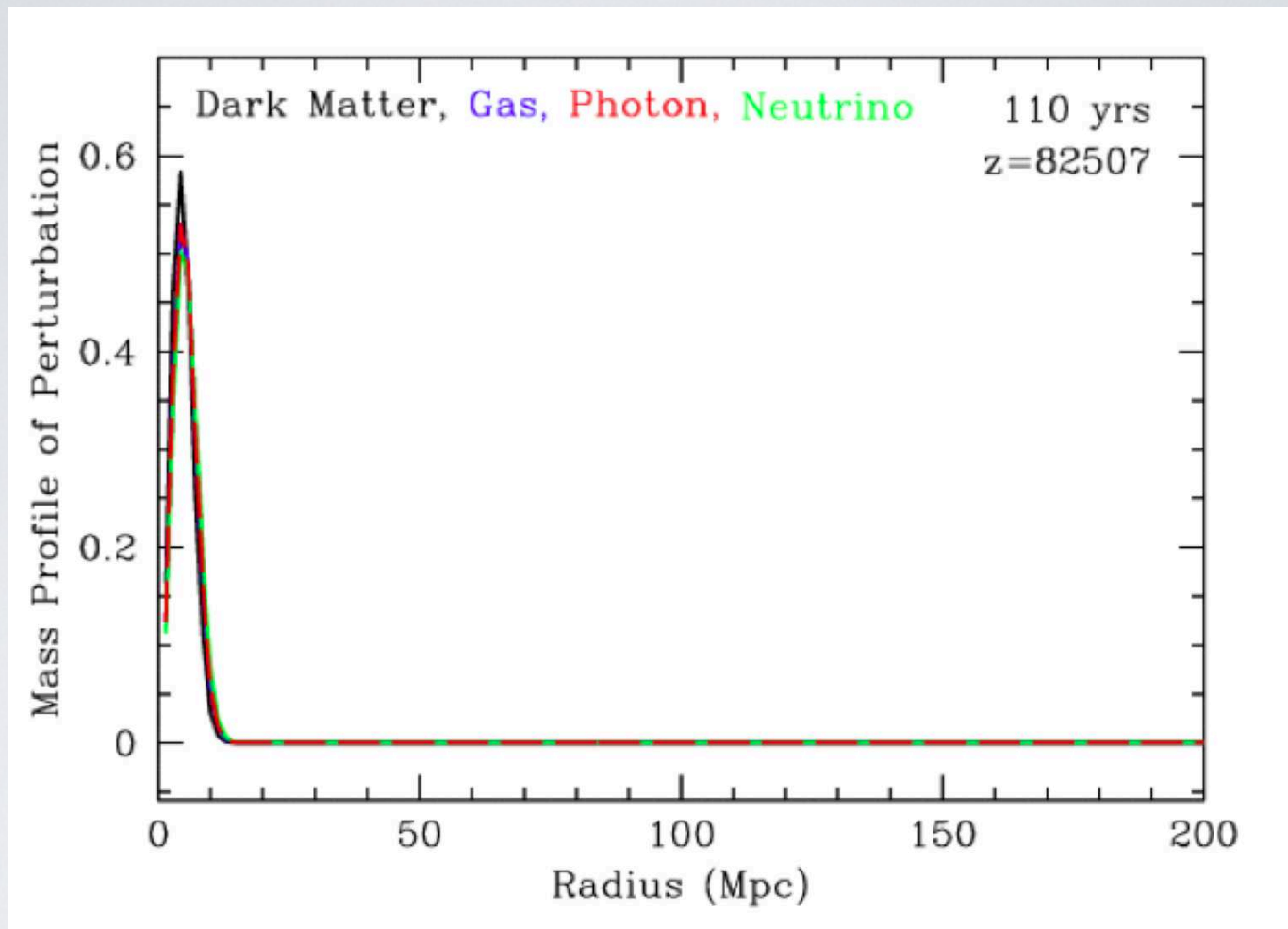


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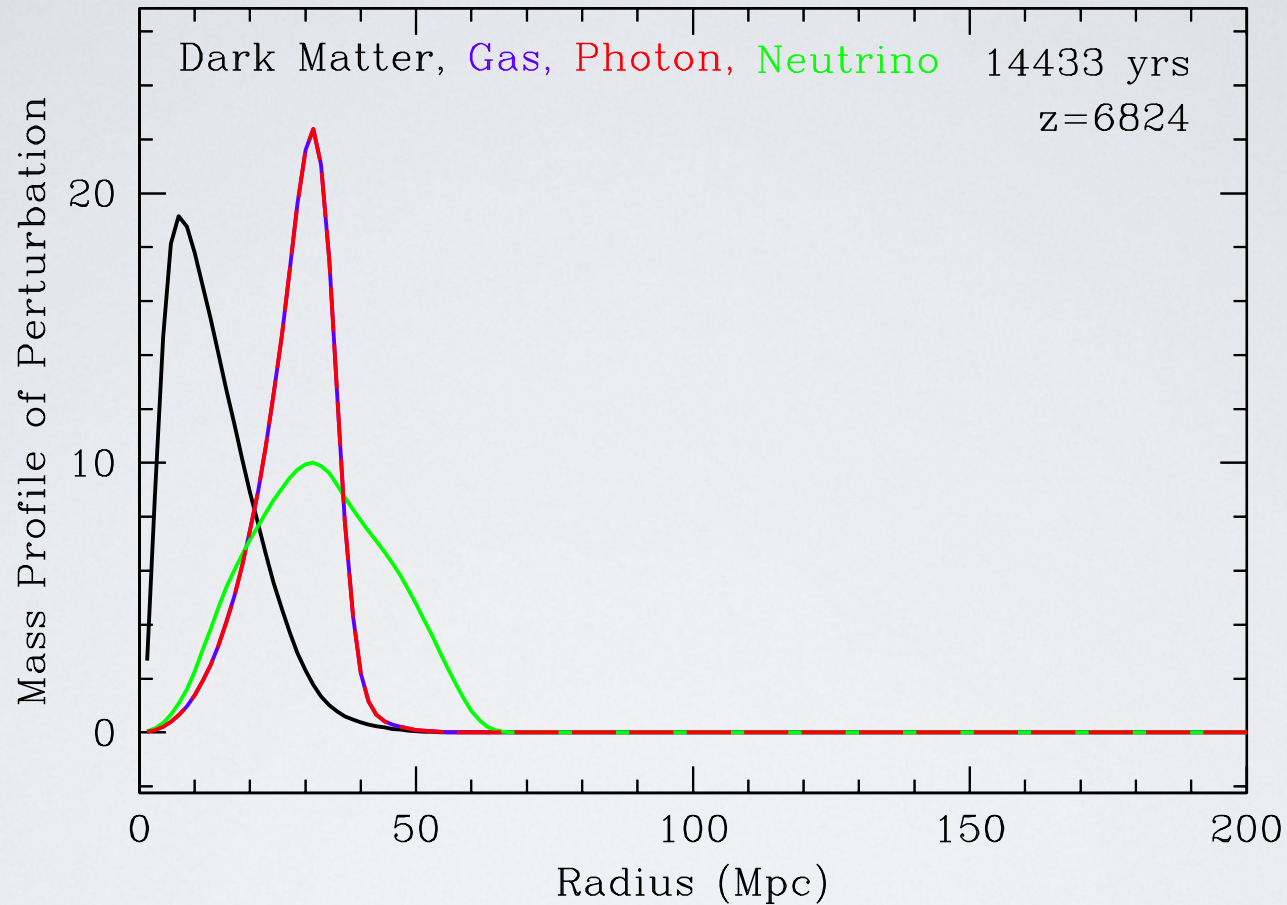
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Distribution of fluids



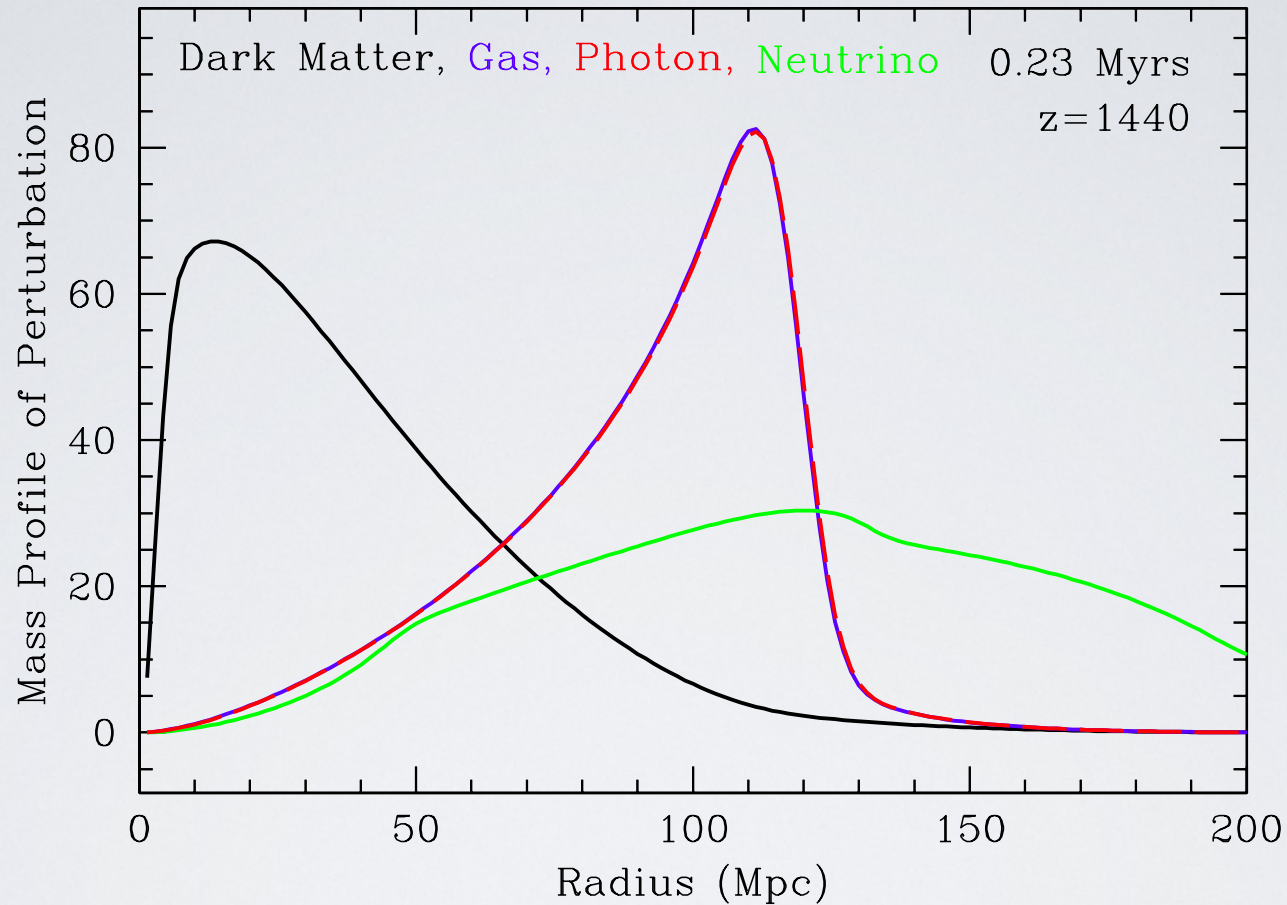
Eisenstein, Seo, and White (2007)

Distribution of fluids



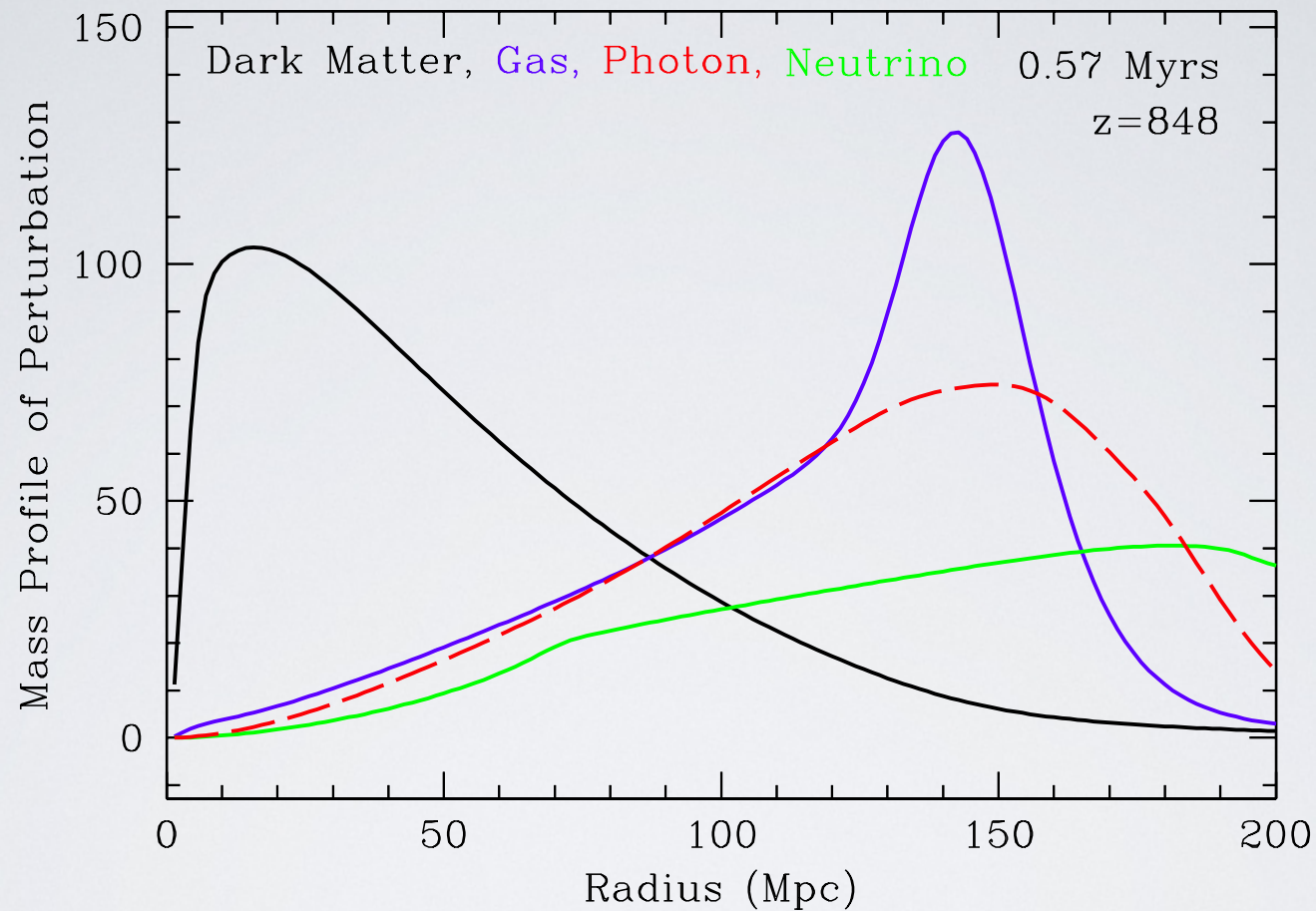
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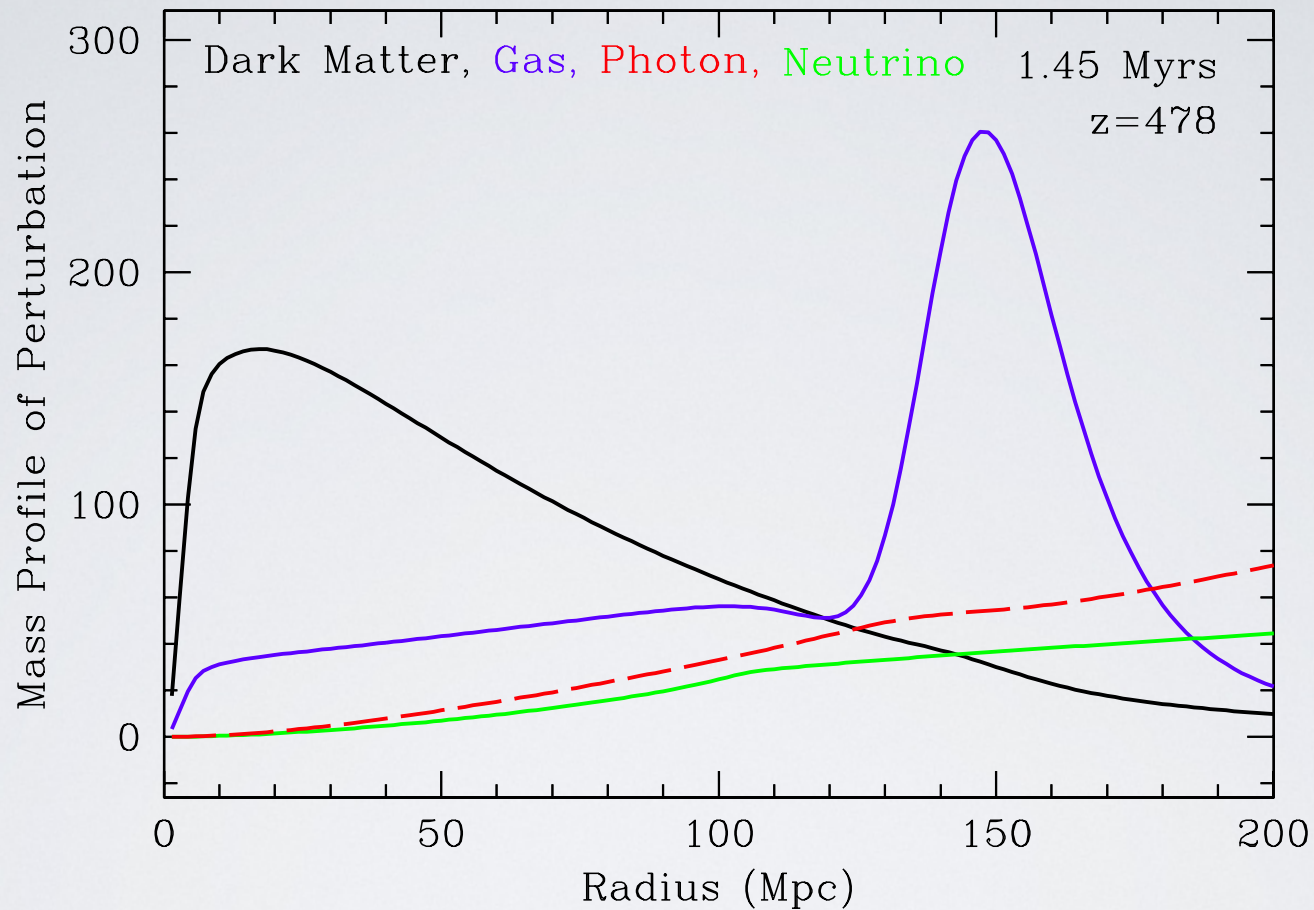
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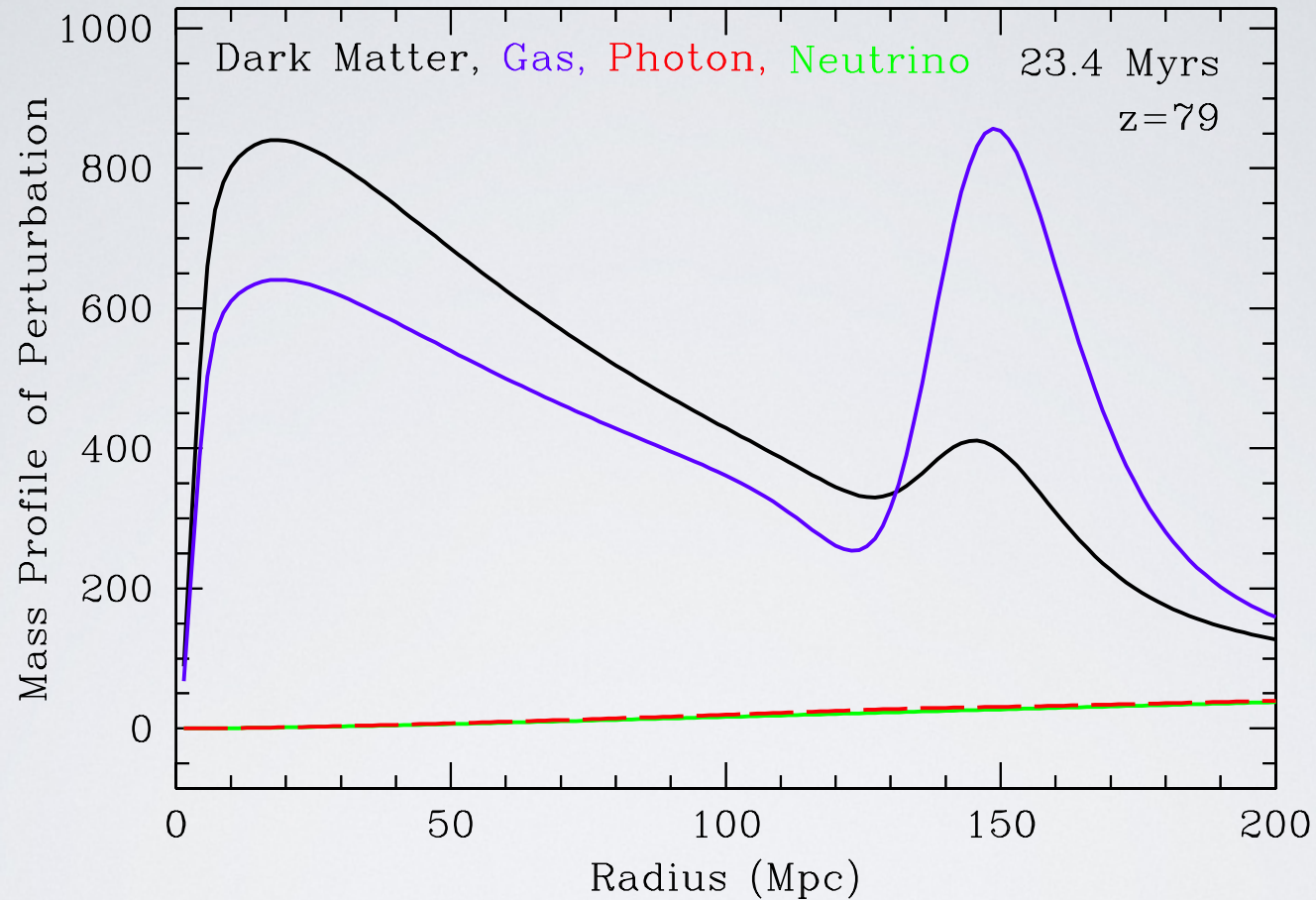
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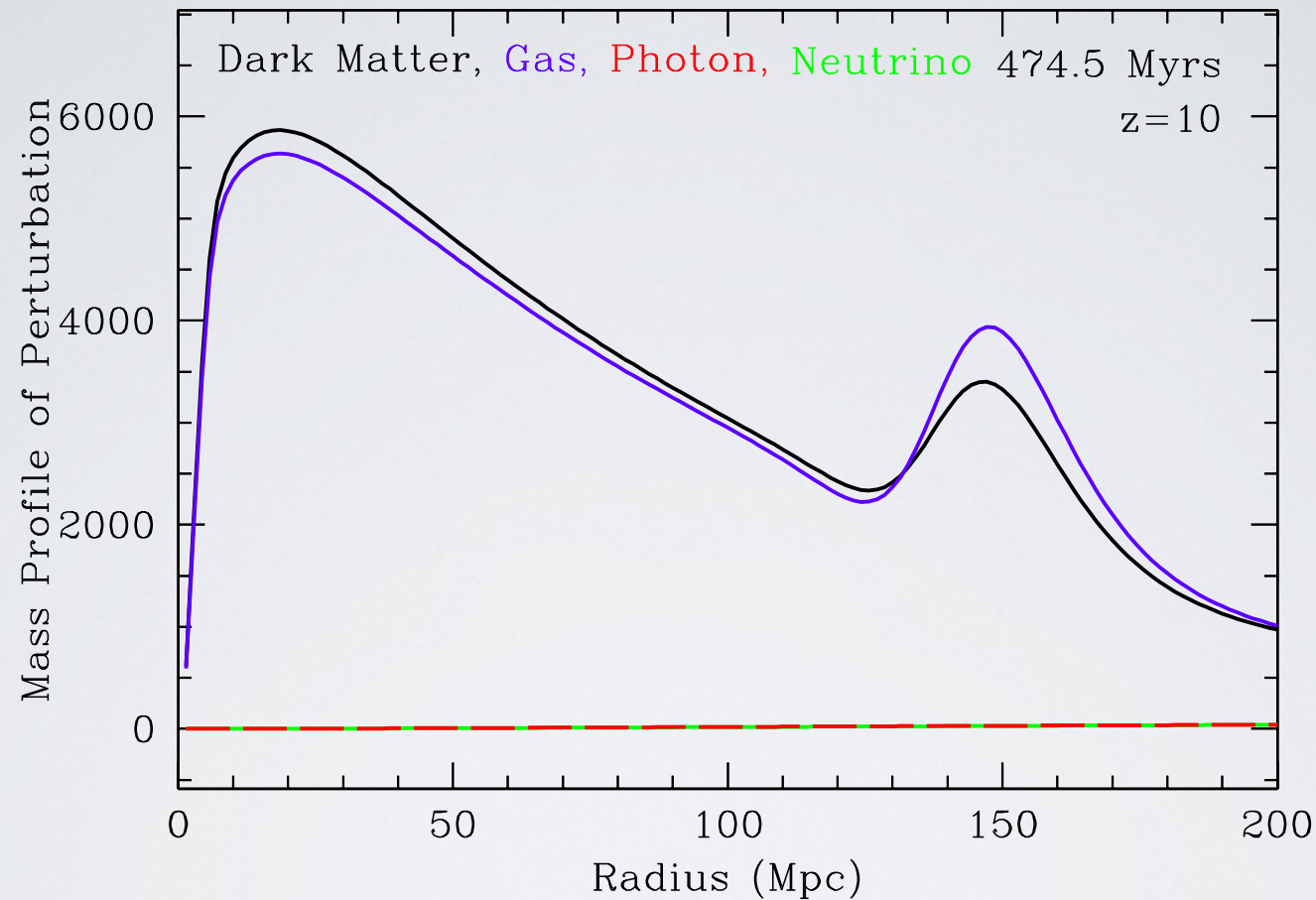
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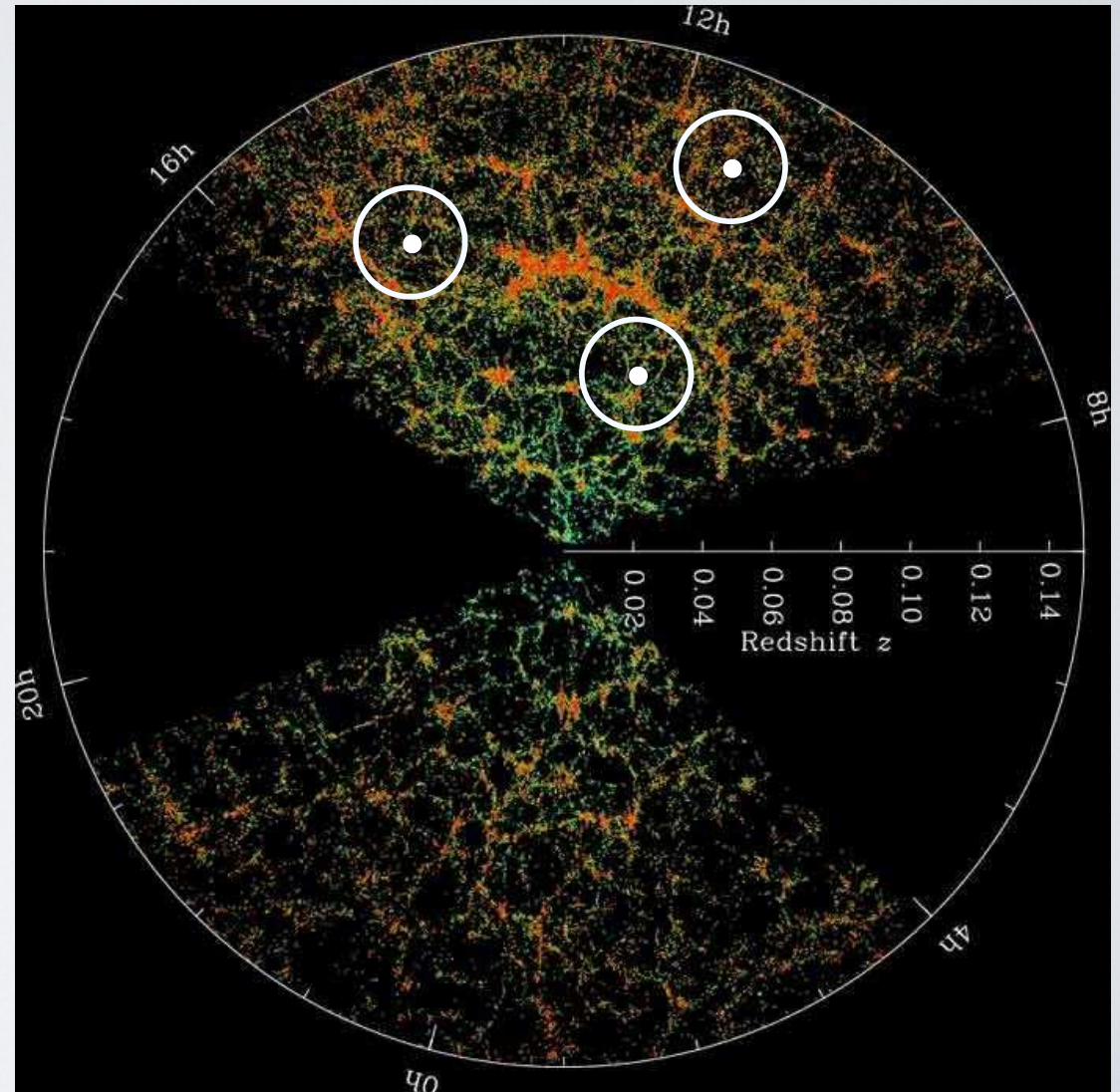
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Galaxy survey

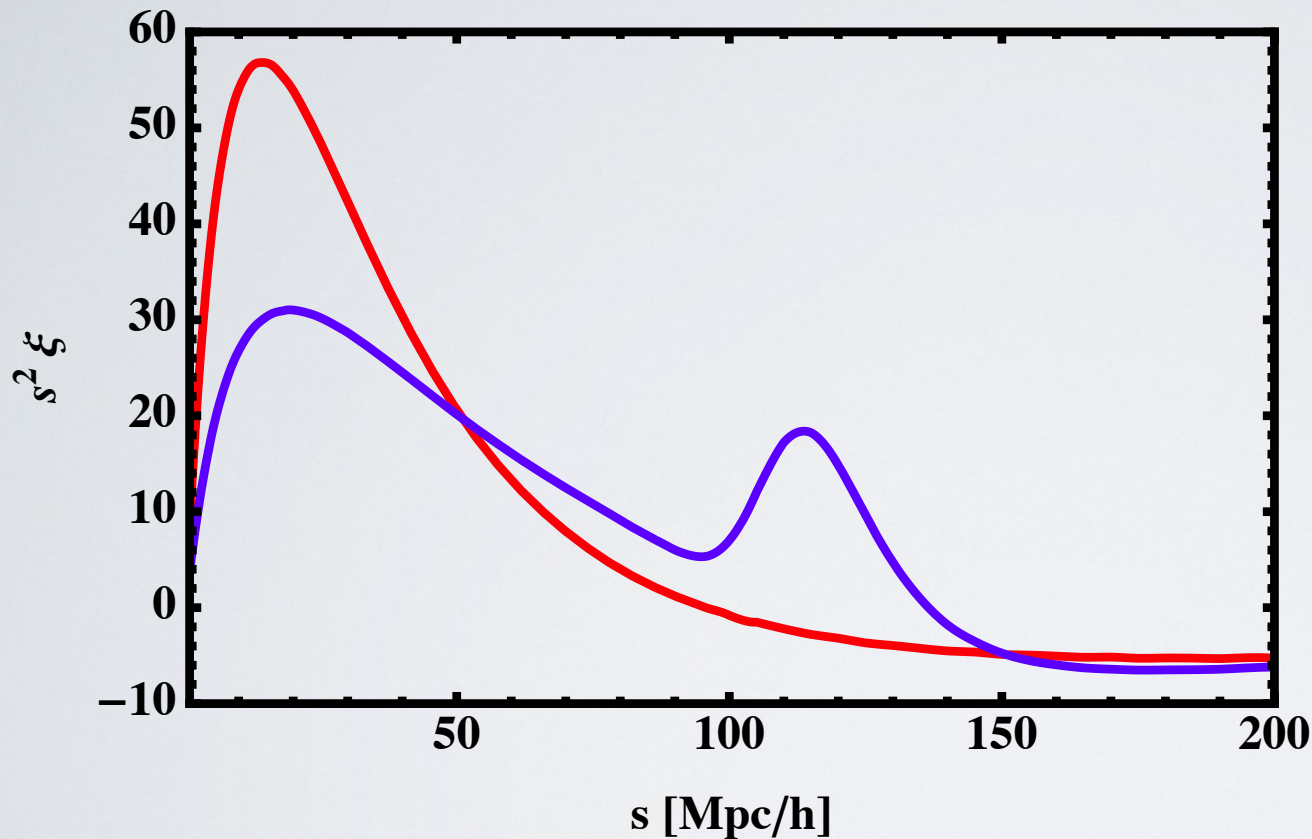
Real universe: superposition of over-densities → **statistical signal**.

- ◆ There is a larger probability to find galaxies separated by 150 Mpc .
- ◆ The two-point correlation function has a **bump** at the sound horizon.

Credit: M. Blanton, SDSS



Correlation function



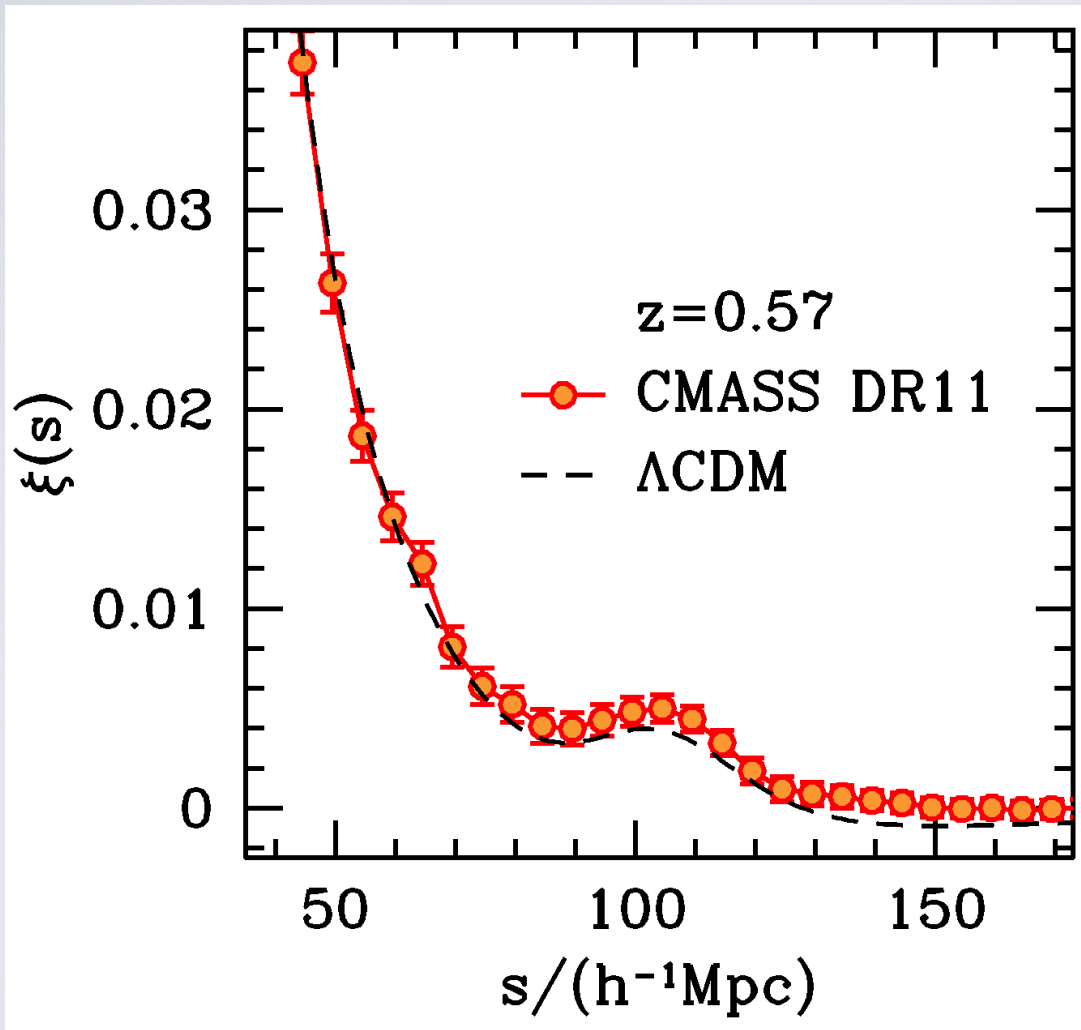
Peak at:

$150 \text{ Mpc} \simeq 110 \text{ Mpc}/h$

Direct **signature** of **baryons**: no bump in a dark matter only universe.

Correlation function

SDSS III, arXiv:1312.4854



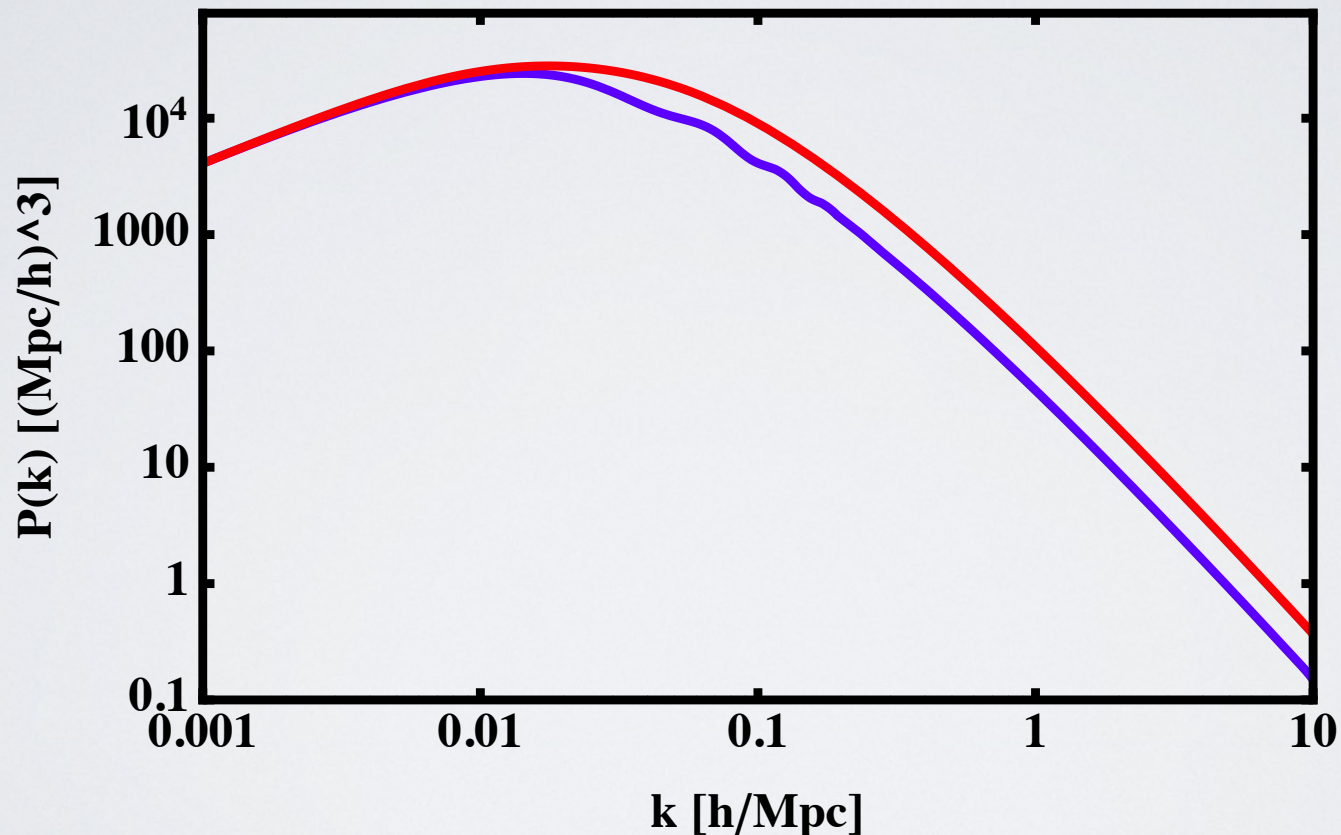
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Power spectrum

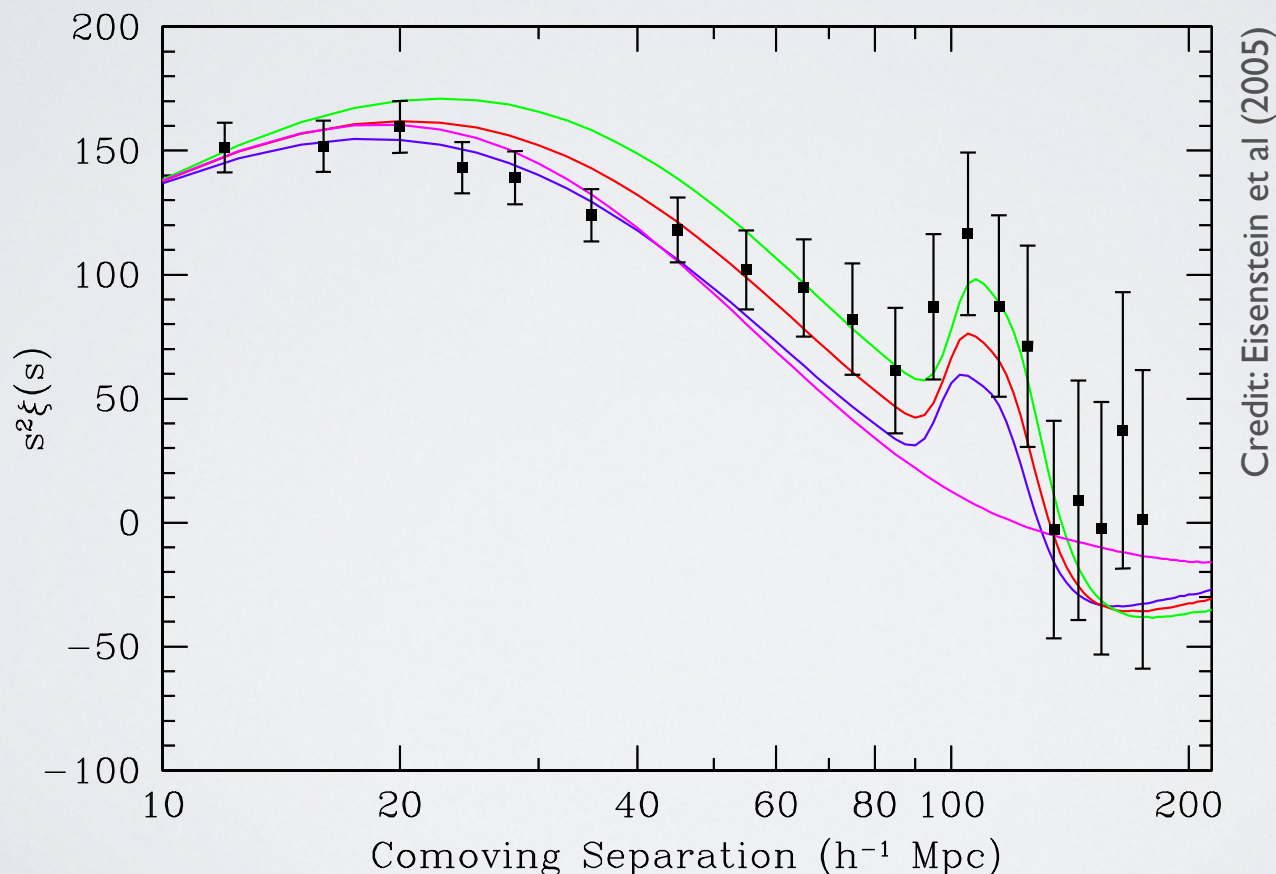
The bump in the correlation function translates into **oscillations** in the power spectrum.



Information

What can we learn from the BAO?

The **amplitude** of the peak depends on the baryons-dark matter ratio



Sound horizon

Position of the peak: distance travelled by the fluid from the big bang until recombination = **sound horizon**.

$$s_h = \int_0^{\eta_{\text{rec}}} d\eta \, c_S = \int_{z_{\text{rec}}}^{\infty} dz \, \frac{c_S}{H(z)}$$

Sound speed:
$$c_S = \frac{1}{3 \left(1 + \frac{3\rho_b}{4\rho_\gamma} \right)}$$

The sound horizon depends on:

- ◆ **Baryon-photon** ratio: more baryon \rightarrow smaller velocity.
- ◆ Time of **decoupling**
- ◆ **Expansion** rate from the big bang until recombination.

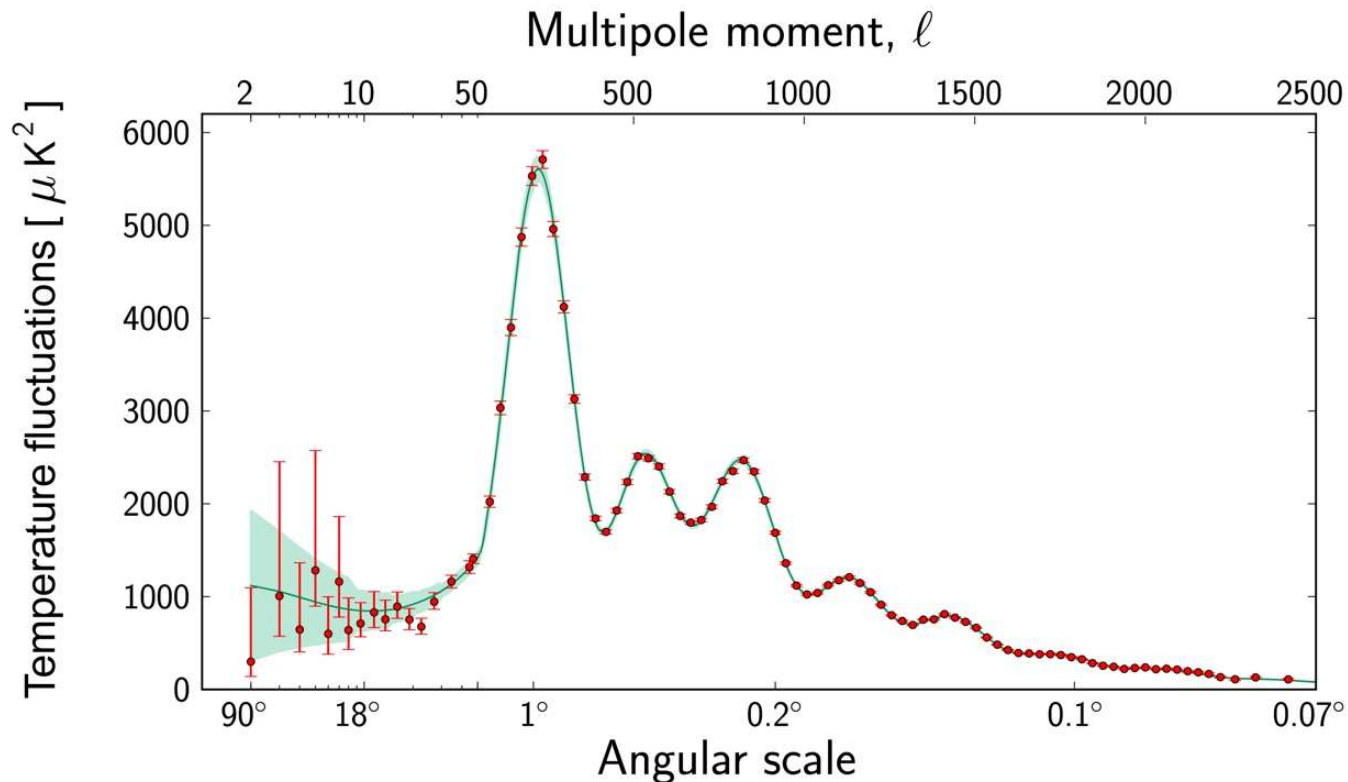
The position of the bump allows to measure these quantities.

Sound horizon in the CMB

The sound horizon can be measured in the CMB.

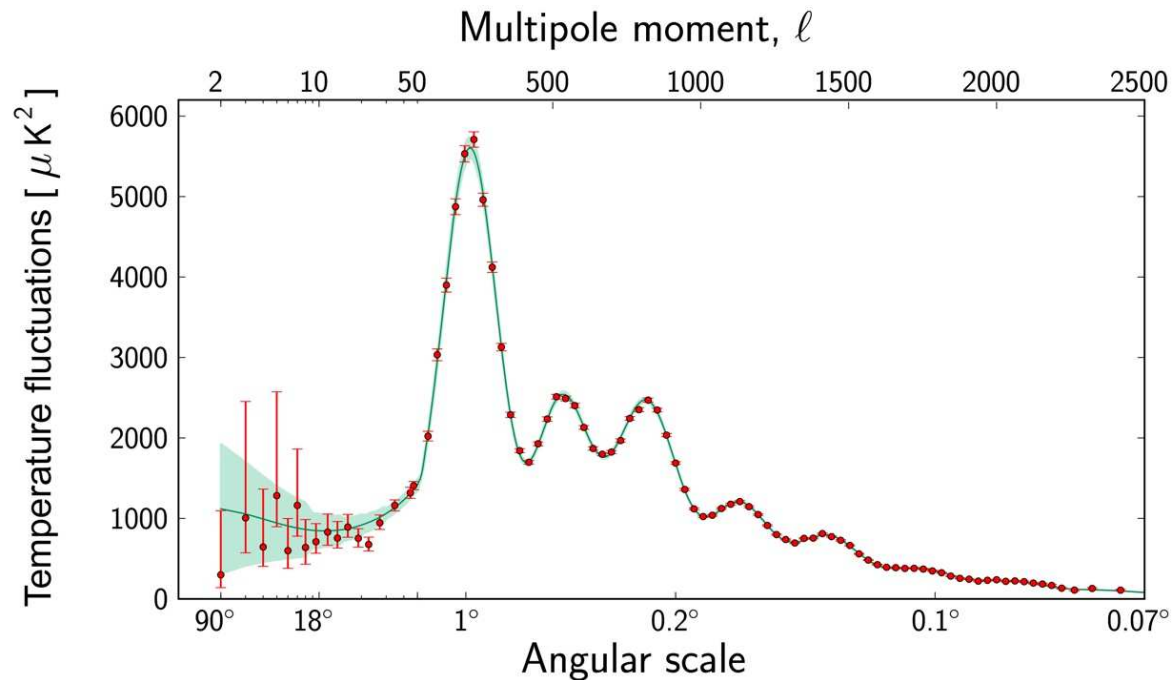
The baryon acoustic **oscillations** are imprinted in the CMB **temperature**.

ESA and the Planck Collaboration



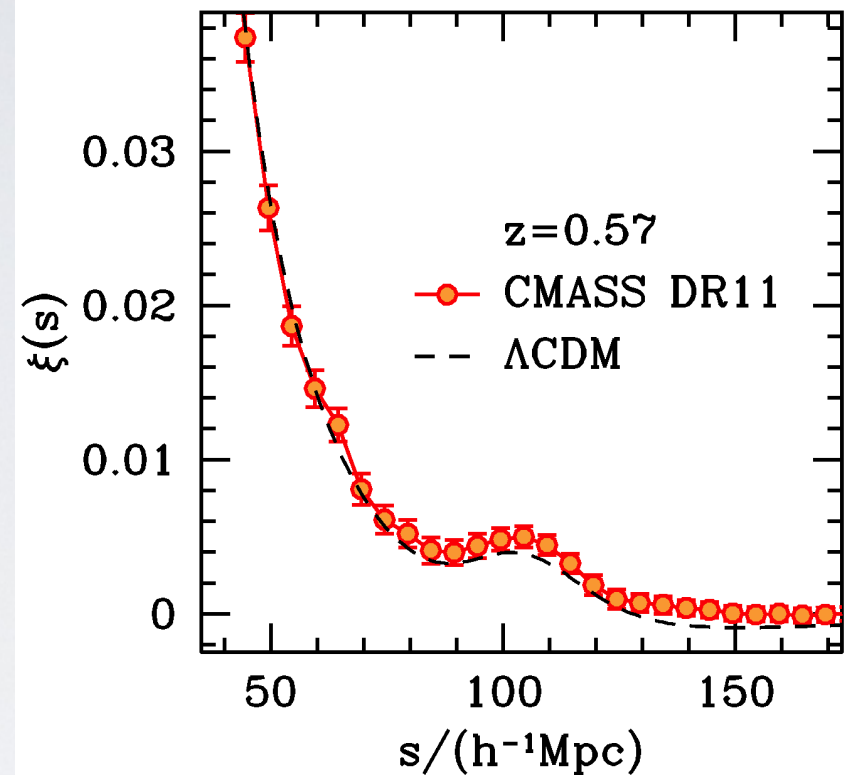
Comparison

CMB



Order one effect

Galaxy correlations



One percent effect

Why is it interesting to measure the BAO scale?

Standard ruler

BAO provide a **standard ruler** → measure the **expansion** rate.

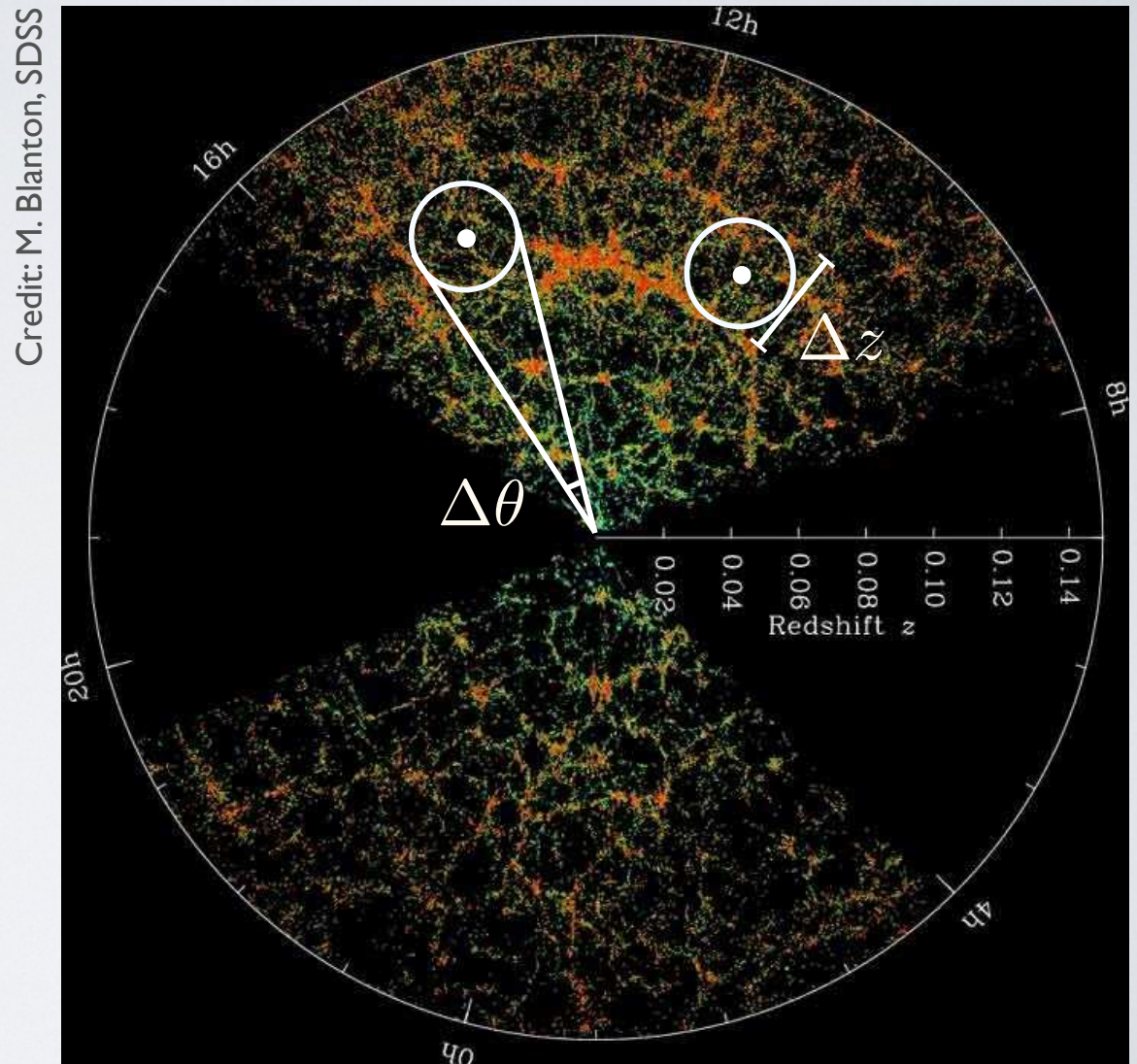
The correlation function peaks at a **comoving** distance of 150 Mpc.

We measure:

- ◆ **angular** separations
- ◆ **redshift** separations

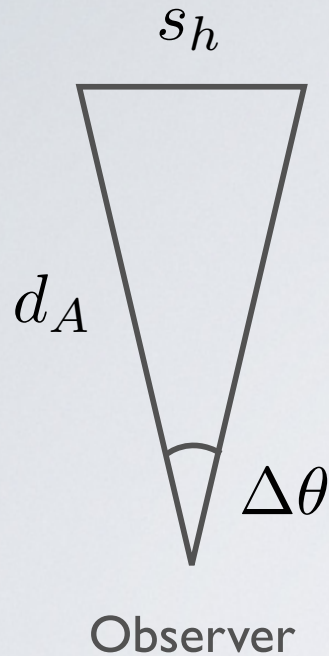
Observed bump at:

$\Delta\theta$ and Δz



Transverse BAO

We can relate $\Delta\theta$ and Δz to s_h (known from CMB)



$$\underset{\substack{\downarrow \\ \text{known}}}{s_h} = d_A \cdot \underset{\substack{\downarrow \\ \text{measured}}}{\Delta\theta}$$

We infer
$$d_A(z) = \frac{1}{1+z} \int_0^z dz' \frac{1}{H(z')}$$

Measuring $\Delta\theta$ provides a measure of the **expansion rate**.

In a flat universe with a cosmological constant:

$$d_A(z) = \frac{1}{1+z} \int_0^z dz' \frac{1}{\sqrt{\Omega_r(1+z')^4 + \Omega_m(1+z')^3 + \Omega_\Lambda}}$$

Luminosity distance

The transverse measurement is very similar to luminosity distance measurements from **supernovae**.

Credit: NASA/JPL-Caltech/R. Hurt (SSC)



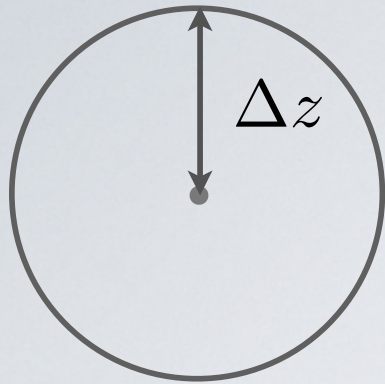
↙

We know the intrinsic **luminosity**, measure the **flux** and infer the distance.

↘

We know the intrinsic **size**, measure the **angular size** and infer the distance.

Radial BAO



•
Observer

Relation between Δz and s_h

$$dt = a dr$$

$$dt = \left(\frac{dt}{da} \right) \frac{da}{dz} dz$$

$$\searrow \quad \frac{1}{\dot{a}} = \frac{a}{H}$$

$$s_h = \frac{1}{H(z)} \Delta z$$

Measure H directly!

Without separation between transverse and radial distance, the BAO peak is measured at the **average distance**

$$D_V(z) = \left[z(1+z)^2 d_A^2(z) H^{-1}(z) \right]^{1/3}$$

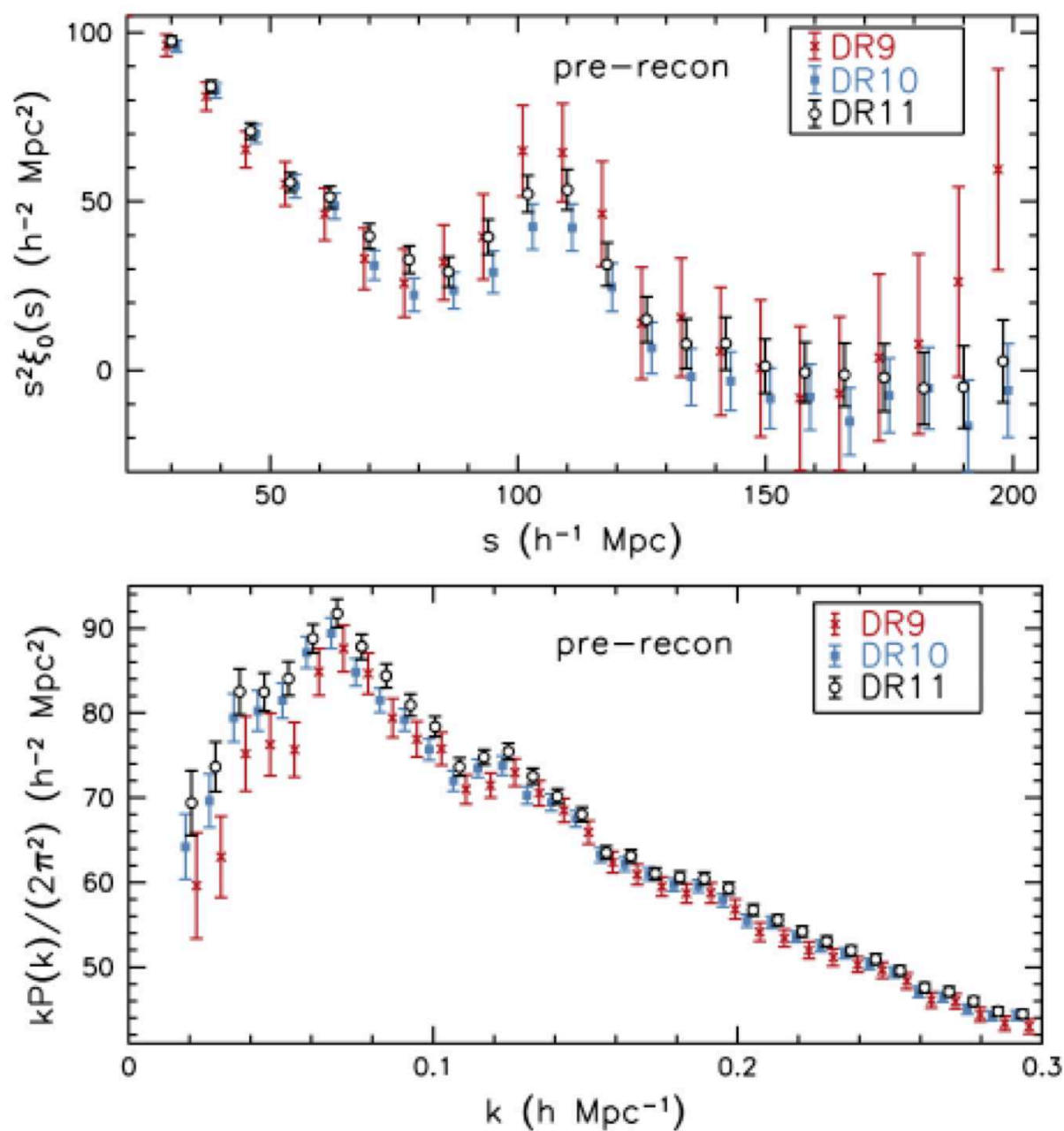
Observations

The second effect of baryons is less noticeable on Figure 7.13 and indeed may never get measured in real life either. [...] It is the traces of these oscillations that are imprinted on the matter transfer function. They are barely (if at all) detectable because baryons are such a small fraction of the total matter. Dodelson: Modern Cosmology p. 209

- ◆ What is the **amplitude** of the BAO excess in terms of the **number** of galaxies? Bassett (2012)
- ◆ Let's pick up one galaxy in the survey and calculate the average **number** of galaxies in a **BAO shell**.
- ◆ Typical number density: $n \sim 10^{-4} \text{ Mpc}^{-3}$
- ◆ Volume of shell: $V = 4\pi \cdot 150^2 \cdot 5 \text{ Mpc}^3 \sim 10^6 \text{ Mpc}^3$
- ◆ Random distribution: 100 galaxies in the shell.

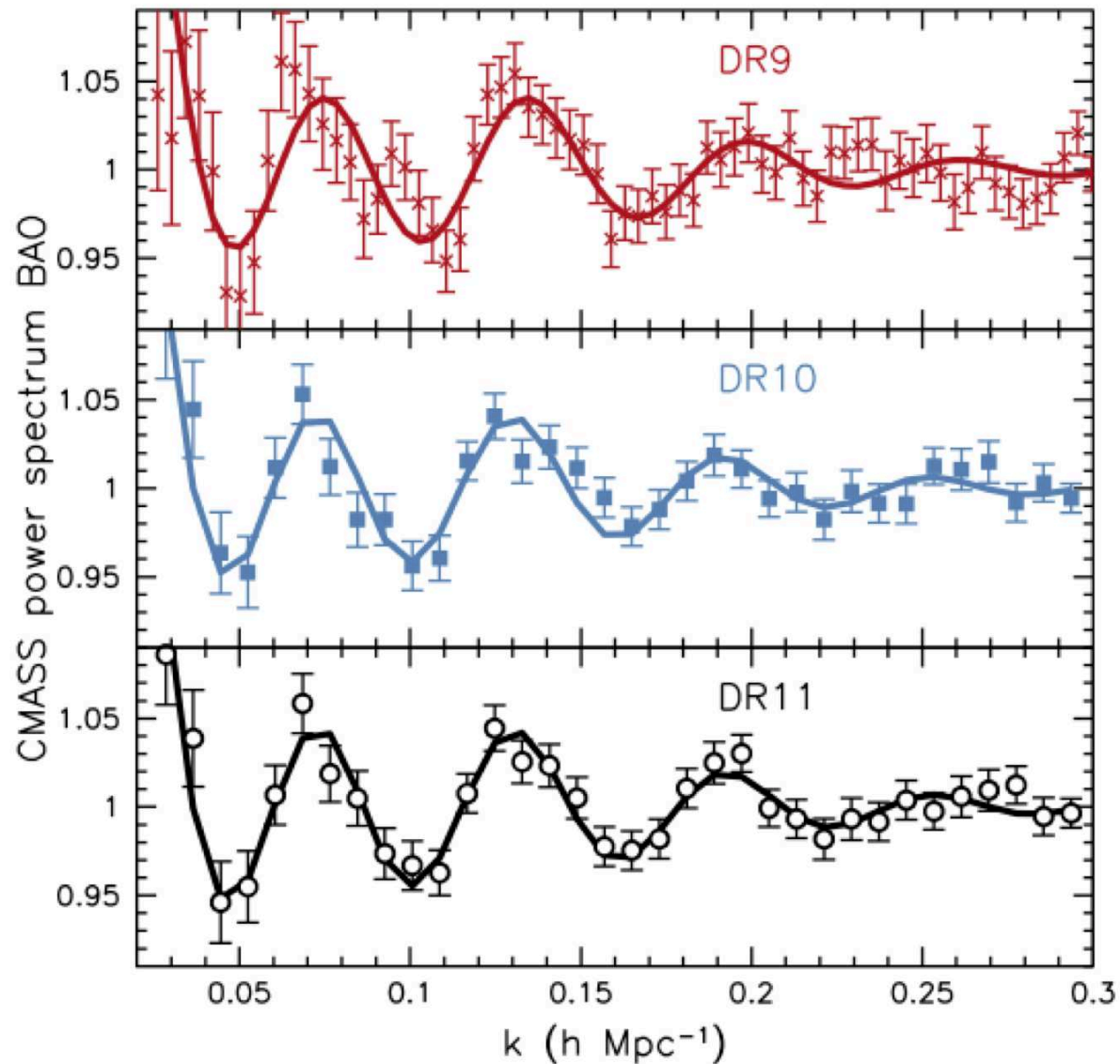
BAO excess 1% \rightarrow **one** extra galaxy!

BOSS results



Credit: Anderson et al. (2013)

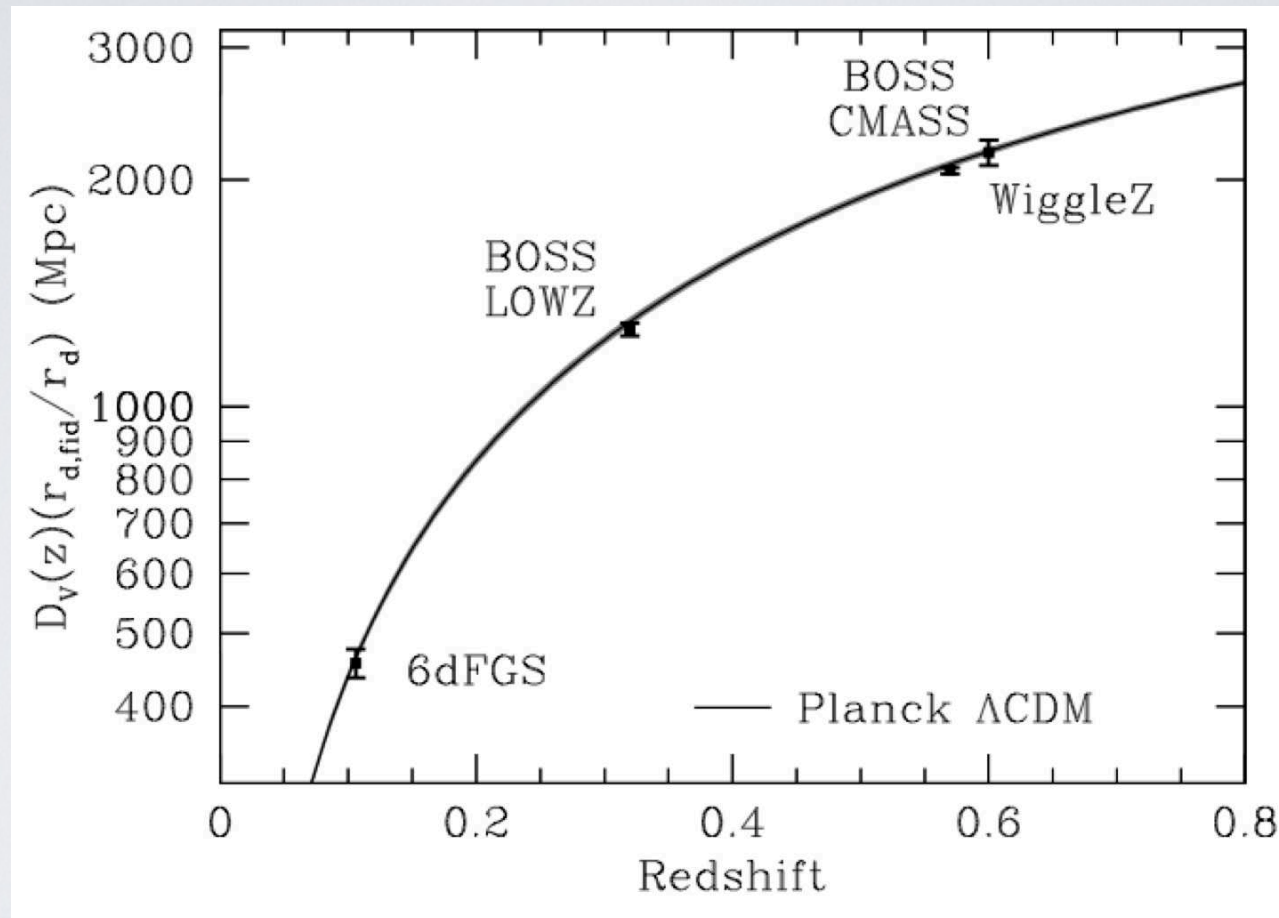
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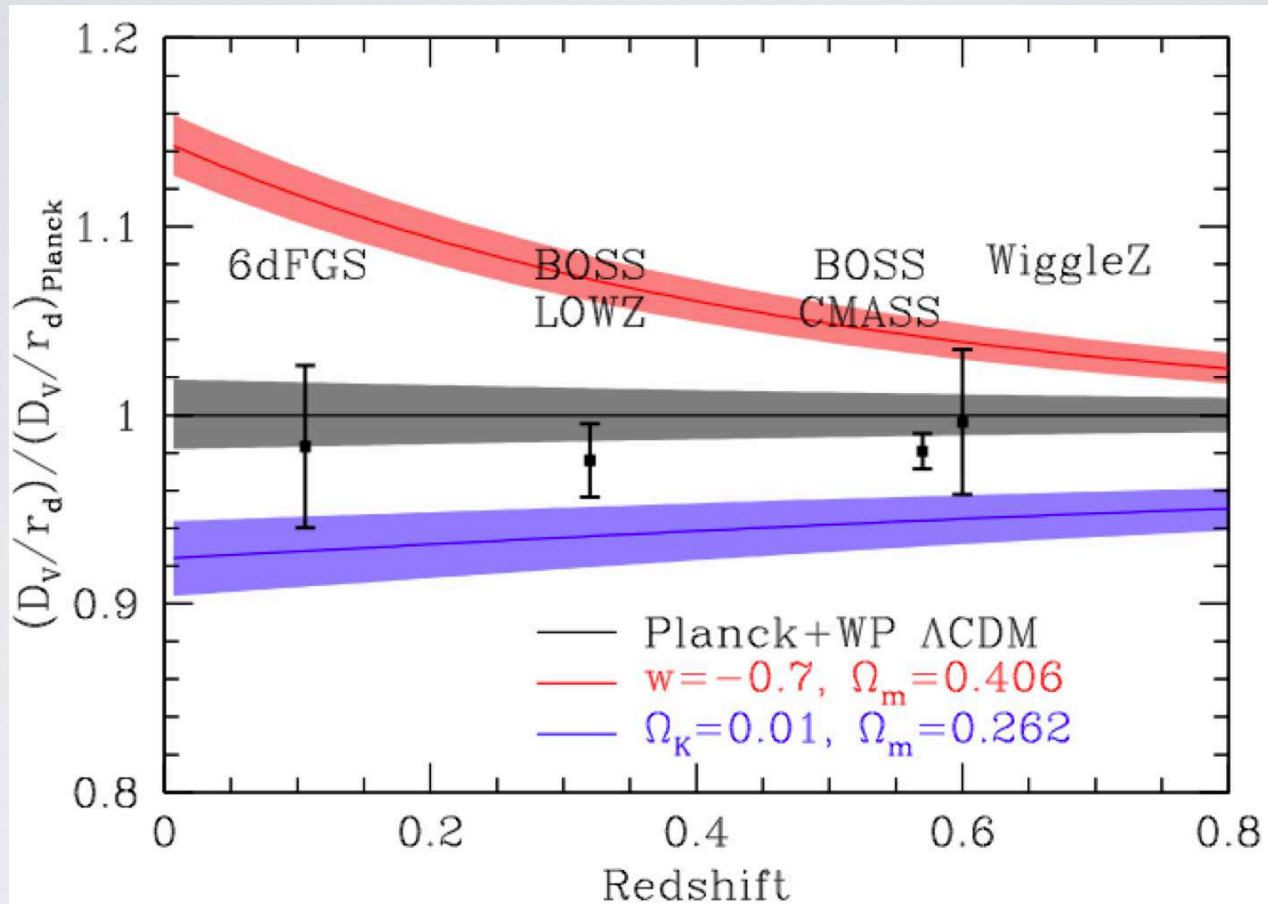
Divided by the smooth spectrum

Distance measurements



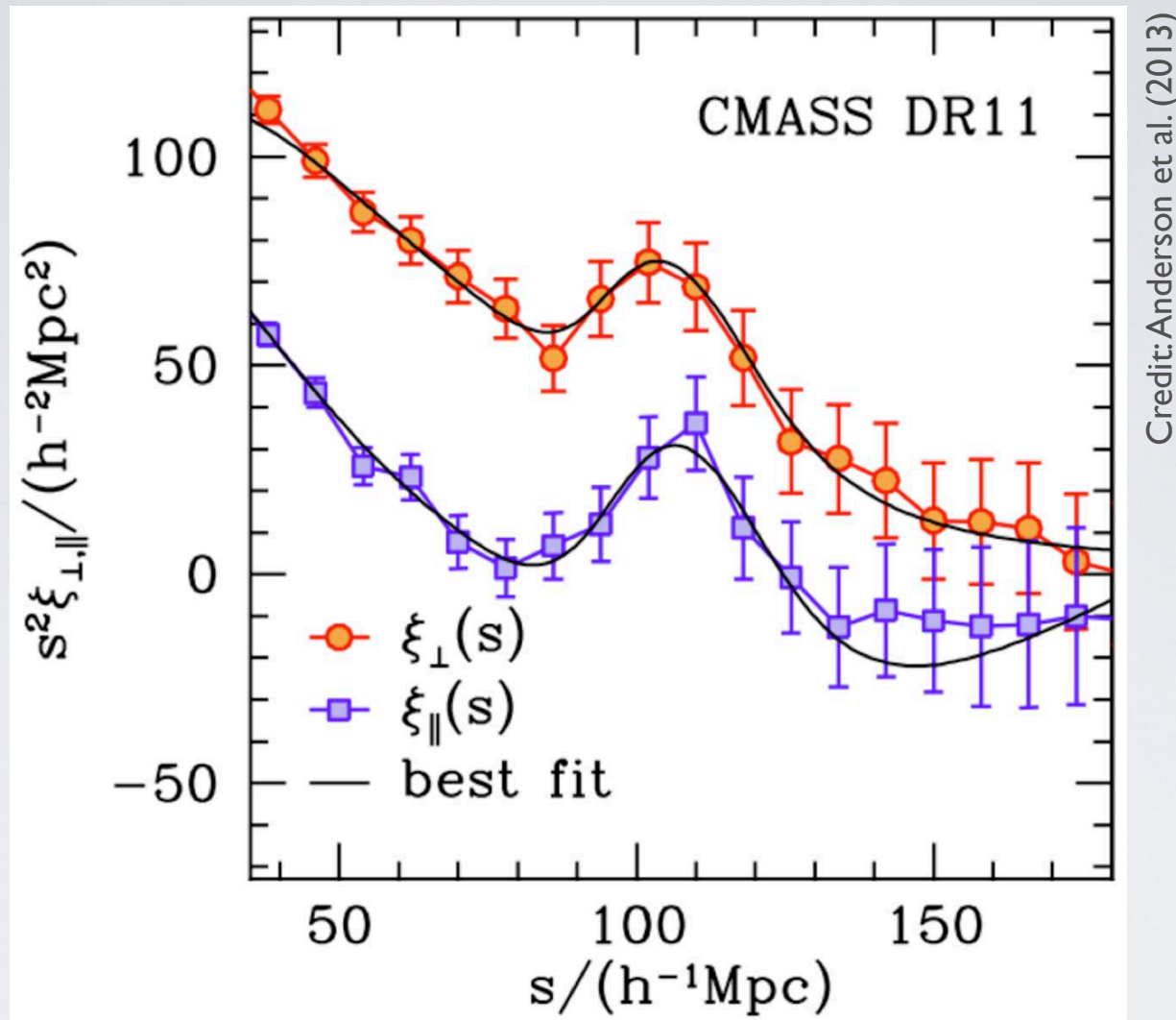
Credit: Anderson et al. (2013)

Accelerating universe



Credit: Anderson et al. (2013)

Separation: radial and transverse



Direct measurement of d_A and H at $z = 0.57$